

# Weak values in the study of quantum correlations

**FOMO**  
Lectures  
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# Objective

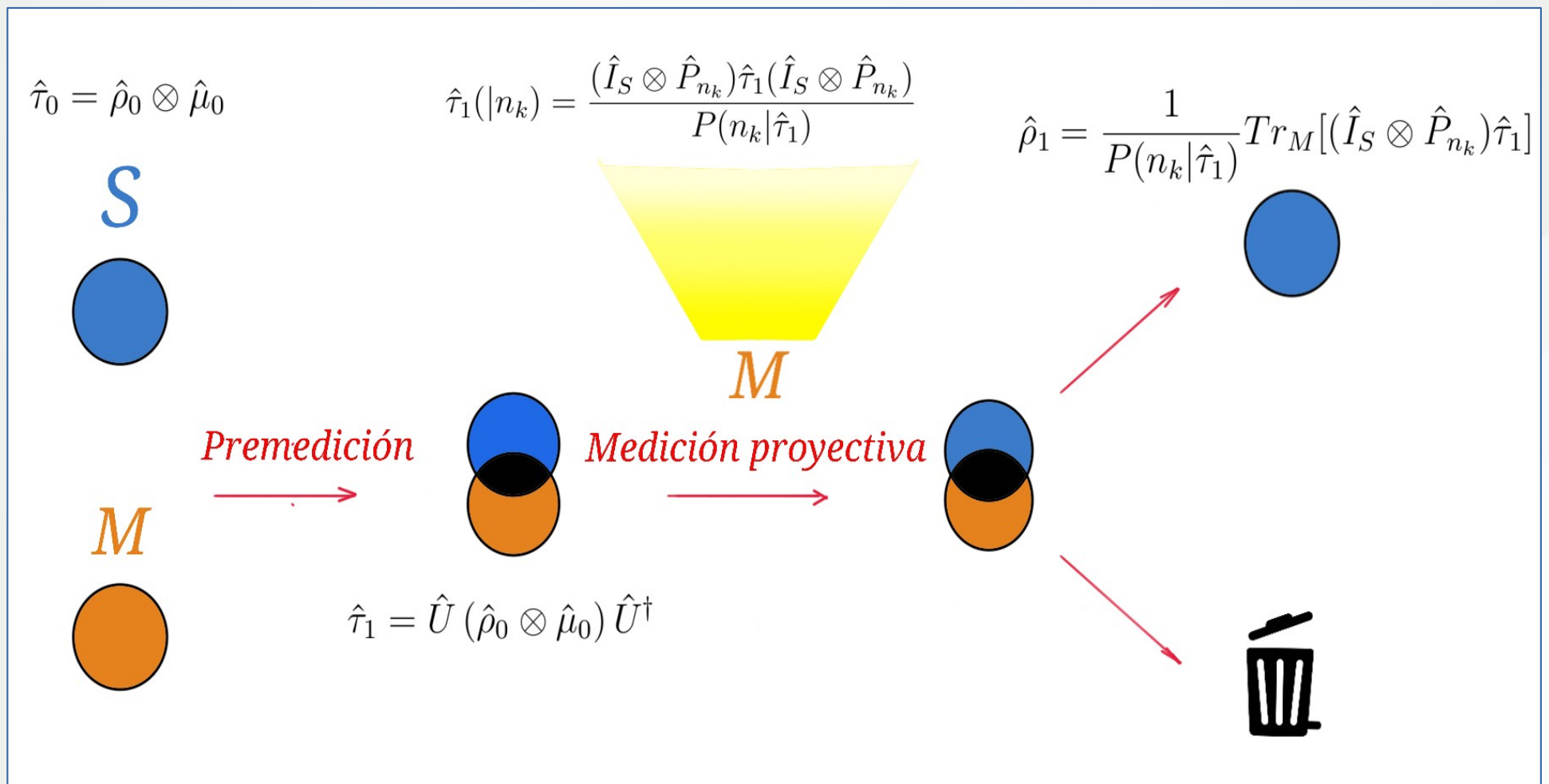
- Determine the utility of weak values to study quantum correlations in general states.

# Structure

1. Weak measurements and weak values
2. Quantum entanglement criterion
3. Further results and Conclusions

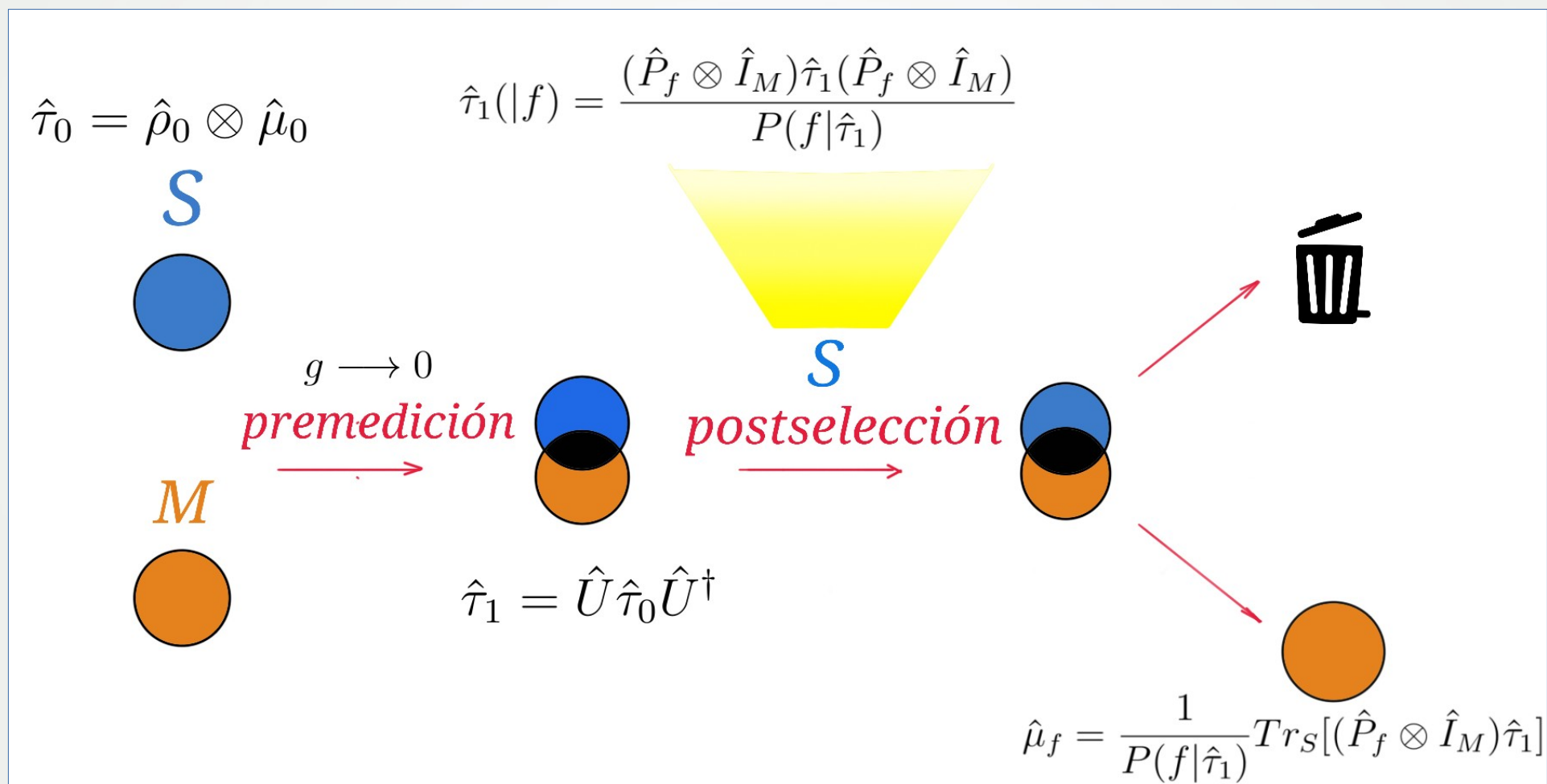
# 1. Weak measurements and weak values

- Ancilla scheme



# 1. Weak measurements and weak values

- Weak measurements with postselection



$$\hat{U} = \exp\left(\frac{-i}{\hbar} g \hat{S} \otimes \hat{N}\right), \quad g = \int_{t-\epsilon}^{t+\epsilon} k(t') dt'$$

# 1. Weak measurements and weak values

- The state of the apparatus after the postselection is

$$\hat{\mu}_f = \hat{\mu}_0 + \frac{2g}{\hbar} \text{Im}(\langle \hat{S} \rangle_w \hat{N} \hat{\mu}_0).$$

With the definition of the weak value for pure states

$$\langle \hat{S} \rangle_w = \frac{\langle f | \hat{S} | s \rangle}{\langle f | s \rangle},$$

## 2. Quantum entanglement criterion


- Entanglement in pure states of bipartite systems (A|B)

$$|\psi\rangle = |\psi\rangle_A \otimes |\psi\rangle_B.$$

In the configuration space

$$\psi(\mathbf{x}_A, \mathbf{x}_B, t) = \psi_A(\mathbf{x}_A, t)\psi_B(\mathbf{x}_B, t),$$

Separability conditions for amplitude and phase


$$\rho(\mathbf{x}_A, \mathbf{x}_B, t) = \rho_A(\mathbf{x}_A, t)\rho_B(\mathbf{x}_B, t),$$

$$S(\mathbf{x}_A, \mathbf{x}_B, t) = S_A(\mathbf{x}_A, t) + S_B(\mathbf{x}_B, t).$$

## 2. Quantum entanglement criterion

- Consider a 2-particles system in the pure state

$$\psi(\mathbf{x}_1, \mathbf{x}_2, t) = \sqrt{\rho(\mathbf{x}_1, \mathbf{x}_2, t)} e^{iS(\mathbf{x}_1, \mathbf{x}_2, t)}.$$

If you calculate

$$\langle \hat{\mathbf{p}}_i \cdot \hat{\mathbf{p}}_j \rangle_w = \frac{\langle \mathbf{x}_1 \mathbf{x}_2 | \mathbf{p}_i \cdot \hat{\mathbf{p}}_j | \psi \rangle}{\langle \mathbf{x}_1 \mathbf{x}_2 | \psi \rangle}$$

$$(i \neq j \in \{1, 2\})$$

$$= \frac{\hat{\mathbf{p}}_i \cdot \hat{\mathbf{p}}_j \psi}{\psi}$$

$$\mathbf{u}_j = \frac{\hbar}{2m_j} \frac{\nabla_j \rho}{\rho} \quad \mathbf{v}_j = \frac{\hbar}{m_j} \nabla_j S$$

$$= \langle \hat{\mathbf{p}}_i \rangle_w \cdot \langle \hat{\mathbf{p}}_j \rangle_w - \hbar m_j (\nabla_i \cdot \mathbf{u}_j + i \nabla_i \cdot \mathbf{v}_j),$$



## 2. Quantum entanglement criterion

If the state is separable, the last term vanishes:

$$\nabla_i \cdot \mathbf{u}_j = 0$$

Sep. in the amplitude

$$\nabla_i \cdot \mathbf{v}_j = 0$$

Sep. in the phase

The criterion can be given in simple terms if one defines the quantity:

Weak correlation

$$C_{\hat{A}\hat{B}}^w = \langle \hat{A}\hat{B} \rangle_w - \langle \hat{A} \rangle_w \langle \hat{B} \rangle_w.$$

In analogy to

$$C_{\hat{A}\hat{B}} = \langle \hat{A}\hat{B} \rangle - \langle \hat{A} \rangle \langle \hat{B} \rangle$$

$$C_{\hat{p}_i \hat{p}_j}^w = -\hbar m_j (\nabla_i \cdot \mathbf{u}_j + i \nabla_i \cdot \mathbf{v}_j)$$

### 3. Further Results and Conclusions

- One can generalize the entanglement criterion to n-particles system in a pure state. What you obtain then is a bipartite entanglement criterion
- You can generalize the weak correlation to mixed states in two non-equivalent ways. The generalizations allow you to establish quantum discord criteria
- One of these generalizations has to do with the skew information defined by Wigner and Yanase
- Weak values are a useful tool for the study of quantum correlations in general systems

Thank you!